

# The hydrodynamics behind the high-energy emission of high-mass binaries with pulsar

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**Variable Galactic Gamma-Ray Sources IV**

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# Gamma-ray High-Mass Pulsar Binaries (HMPB)

- The colliding wind region is typically considered the high-energy emitter: the shocked and the unshocked pulsar wind, the shocked stellar wind, and electromagnetic cascades.

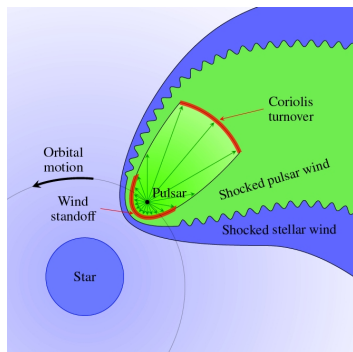
(e.g. Maraschi & Treves 1981; Tavani & Arons 1997; Kirk et al. 1999; Sierpowska & Bednarek 2005; Dubus 2006; Khangulyan et al. 2007; Neronov and Chernyakova 2007; Derishev & Aharonian 2012)

- In this talk we assess the potentiality of middle scales as the high-energy emitter.

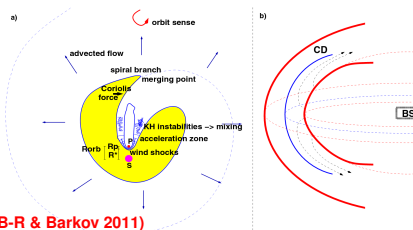
(see B-R & Barkov 2011; Bednarek & Sitarek 2013)

- The flow on middle scales is considered here to be well described by relativistic hydrodynamics (i.e. ideal RMHD;  $\lambda \ll R_{\text{orb}}$ ; low- $B$ ).

(see Bogovalov et al. 2012; see discussion in B-R et al. 2015)



(Zabalza et al. 2013)



(B-R & Barkov 2011)

1 Introduction

**2 Binary scales**

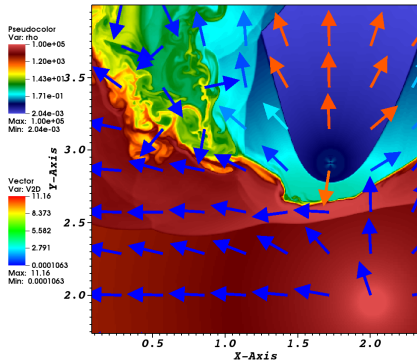
3 Middle scales

4 High-energy emission on middle scales

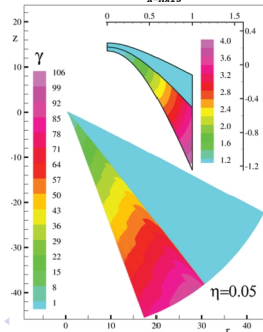
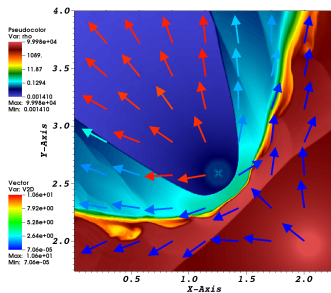
# The colliding wind region

- In HMPB, the massive star and the non-accreting pulsar winds collide.
- After being shocked, the pulsar wind quickly reaccelerates while instabilities develop in the CD.

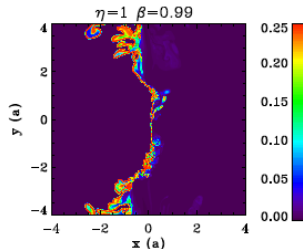
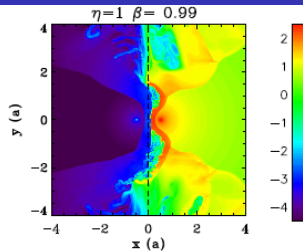
(B-R, Barkov, Khangulyan & Perucho 2012) ↓ ↑



(Bogovalov et al. 2008) →



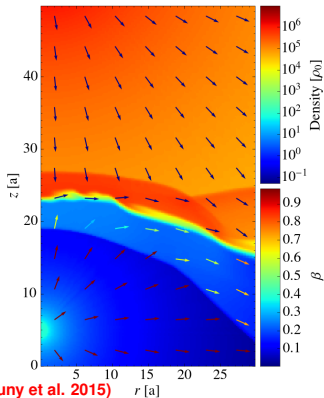
# Instabilities within the colliding wind region



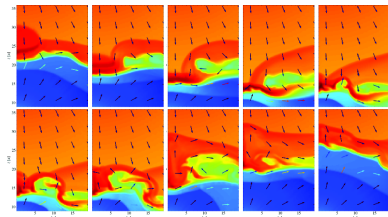
(Lamberts et al. 2013)

Analytic works on the shocked flow unstable nature:

e.g. B-R & Barkov (2011); Bednarek & Sitarek (2013)



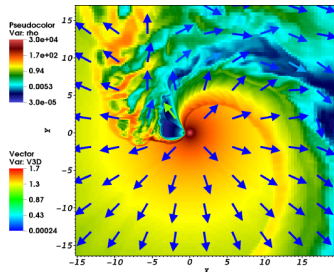
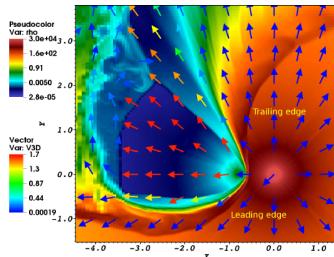
(Paredes-Fortuny et al. 2015)



# Effects of the orbit on binary scales and beyond

- The stellar wind exerts a strong Coriolis force on the pulsar wind due to orbital motion.
- The reaccelerated and the unperturbed pulsar winds are terminated away from the star.
- Orbital motion deflects the shocked winds; instabilities from the CD:
  - \*Rayleigh-Taylor (RT)
  - \*Richtmyer-Meshkov (impulsive RT)
  - \*Kelvin-Helmholtz
- The shocked pulsar wind suffers reacceleration, turbulence, further shocks, and stellar wind entrainment.

(B-R, Barkov & Perucho 2015)

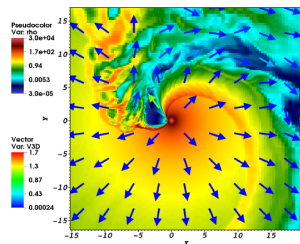
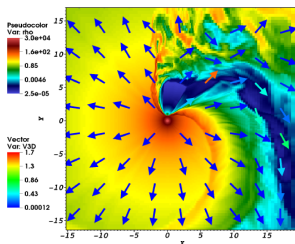
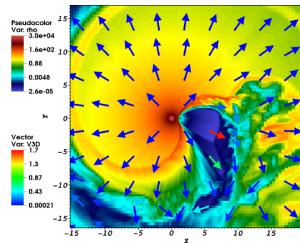
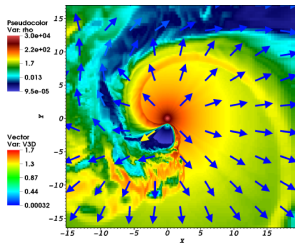




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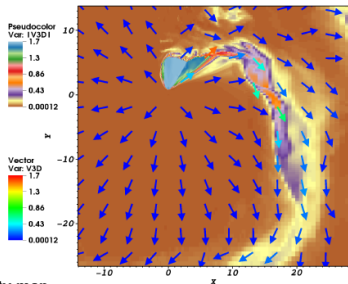
# Middle scales: low eccentricity case (I)

- A one-arm spiral forms made of shocked stellar wind channelling shocked pulsar wind.
- Strong turbulence and mass entrainment from the CD occur.
- The shocked pulsar wind is slowed down by mass-load, shocks and turbulence.
- Pulsar wind energy is accumulated and smoothly evolves along the orbit.

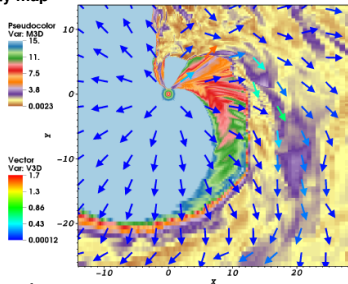


(B-R, Barkov & Perucho 2015)

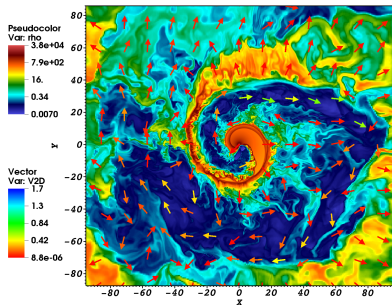
# Middle scales: low eccentricity case (II)



Velocity map



Mach number map



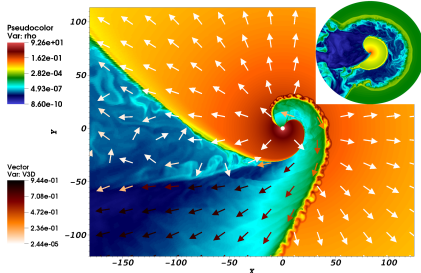
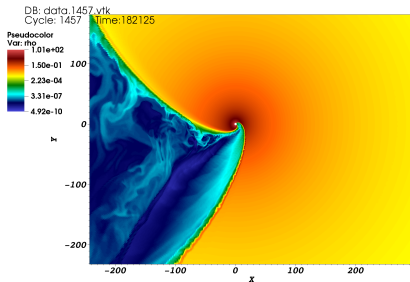
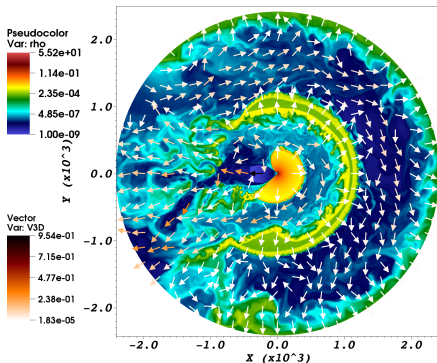
Larger scale density map

- The shocked pulsar wind is slow and sub- or trans-sonic.
- Shocked pulsar wind expansion and instability growth lead to spiral merging and disruption after one turn.

(B-R, Barkov, Khangulyan & Perucho 2012, B-R, Barkov & Perucho 2015 for LS 5039;  $e \approx 0.24$ )

# Middle scales: high eccentricity case (I)

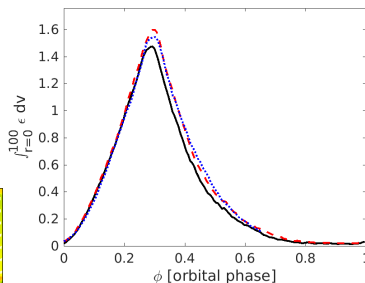
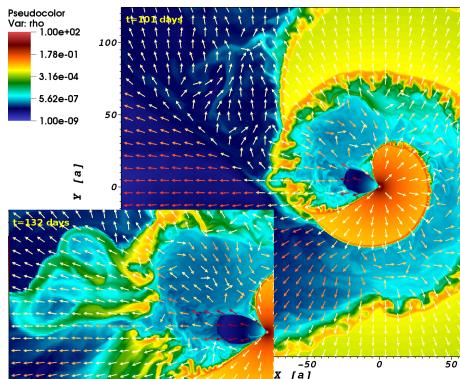
- If eccentricity is high, the pulsar wind is mostly directed apastron-wards.
- The flow accumulates from periastron to apastron, and then is released due to spiral disruption at the apastron side.



(Barkov & B-R 2016 for PSR B1259-63;  $e = 0.87$ )

# Middle scales: high eccentricity case (II)

- The impact of the pulsar wind blows away the shocked stellar wind close to apastron.
- The material blown away may be detectable on large scales.



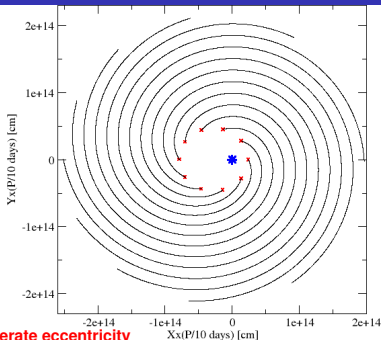
- The pulsar wind energy accumulates in the first spiral turn and peaks around apastron,
- Then, the energy is quickly advected away.

# Outline

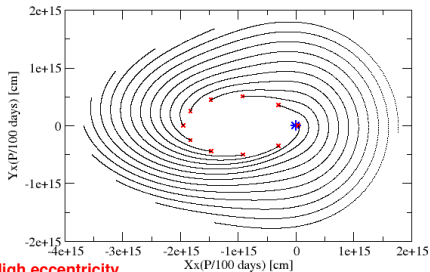
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# Shocked wind spiral

- Initially, the bent shocked flow follows a quasi-ballistic trajectory dominated by the large stellar wind mass.
- The shocked pulsar wind path is determined by the quasi-ballistic trajectory of the stellar wind.
- The spiral shape is characterized by starting at  $r \sim 3 v_w / \Omega_{\text{orb}}$  and  $\phi$  from periastron, and the previously injected matter is radially displaced by  $v_w$ .
- As pulsar wind energy accumulates in the region, and instabilities develop, the ballistic trajectory is lost, and the spiral tends to disrupt.



Moderate eccentricity



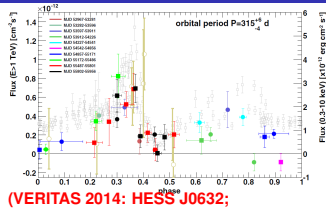
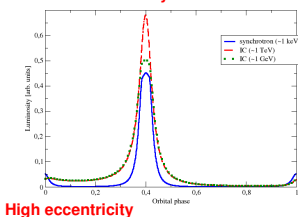
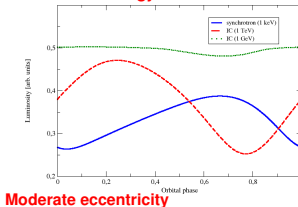
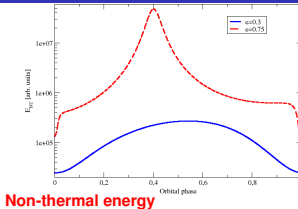
High eccentricity

(B-R & Barkov, in prep.)

# Non-thermal emission

- Assuming particle injection at the spiral onset, and the flow velocity, the accumulated energy can be computed.
- Radiation efficiencies are high  $\lesssim 100\%$ .
- Approximate synchrotron and IC lightcurves can be derived.

(B-R & Barkov, in prep.)



see also LS I +61 303)

