



# The influence of the stellar wind on the jets of high-mass microquasars

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### Physical scenario

- High-mass microquasar in which a strong stellar wind interacts with the jets in a number of ways:
  - Recollimation shock
  - Bending
  - Large scale helical structure



• Numerical and analytical studies predict the formation of a recollimation shock at the binary scales



3D relativistic simulations by Perucho et al. 2010

#### • For low jet powers, jets may be disrupted at very small scales



• No recollimation for high jet powers:

Bosch-Ramon & Barkov 2016

$$L_{\rm j} \gtrsim 2.4 \times 10^{37} \left( \frac{\dot{M}_{\rm w}}{10^{-6} {\rm M}_{\odot} {\rm yr}^{-1}} \right) \left( \frac{v_{\rm w}}{2 \times 10^8 {\rm cm s}^{-1}} \right) \frac{\gamma_{\rm j} (\gamma_{\rm j} - 1)}{\beta_{\rm j}} {\rm ~erg~s}^{-1}$$



Yoon et al. 2016

• Jets are also bent due to the wind impact on them. For small enough bending angles ( $\phi \leq 30 \text{ deg}$ ): Bosch-Ramon & Barkov 2016

$$\phi \approx 17 \left(\frac{L_{\rm j}}{10^{37} {\rm erg \, s^{-1}}}\right)^{-1} \left(\frac{\theta_{\rm j}}{0.1 \ {\rm rad}}\right) \left(\frac{\dot{M}_{\rm w}}{10^{-6} {\rm M}_{\odot} \ {\rm yr^{-1}}}\right) \left(\frac{\nu_{\rm w}}{2 \times 10^8 {\rm cm \, s^{-1}}}\right) \frac{(\gamma_{\rm j} - 1)}{\gamma_{\rm j} \beta_{\rm j}} \ {\rm deg}$$

- Bending combined with orbital motion could lead to a helical pattern if the jet is not disrupted before and  $\phi > \theta_{\rm j}$
- Significant mixing of wind and jet material is expected already within the binary scales

Perucho et al. 2010, 2012



#### Cygnus X-3



Mioduszewski et al. 2001



Miller-Jones et al. 2004

#### Cygnus X-1?



Stirling et al. 2001

### Helical jet model: dynamics

- Jet trajectory is computed from momentum transfer by an isotropic stellar wind
- Within the binary system scales, the jet is bent away from the normal to the orbital plane
- At larger scales, orbital motion makes the jet move towards a helix-like trajectory

Mass-loss rate	$\dot{M}_{ m w}$	$10^{-6}~{ m M}_{\odot}~{ m yr}^{-1}$
Terminal wind speed	$\mathcal{V}_{\infty}$	$2 \times 10^8 \text{ cm s}^{-1}$
Jet luminosity	$L_{j}$	$5 \times 10^{36} \text{ erg s}^{-1}$
Orbital separation	a	$3 \times 10^{12} \text{ cm}$



- Nonthermal electrons are injected at the recollimation shock with a power of  $0.1L_j$
- Particle distribution computed at each point along the jet
- Cooling:
  - Adiabatic
  - Synchrotron
  - Inverse Compton
- Emission:
  - Synchrotron
  - Inverse Compton
- Absorption:
  - Free-free absorption by wind ions (radio)
  - Gamma-gamma absorption by stellar photons (VHE  $\gamma$ -rays)

- For higher  $\gamma_j$ , the change in the Doppler boosting does not compensate for the increased energy losses, except for very small inclinations
- The helical structure enhances absorption and affects the IC emission





Molina et al. 2019

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#### Simulated radio maps at 5 GHz



Molina & Bosch-Ramon 2018

#### Simulated radio maps at 5 GHz



#### Molina & Bosch-Ramon 2018

#### Summary

- Jet-wind interaction must be considered when studying HMMQ jets
- Dynamical effects:
  - Jet bending
  - Recollimation shock
  - Jet disruption
  - Helical structure at larger scales
- Non-thermal radiation:
  - Angle dependent quantities affected by the presence of a helical jet structure (IC, γγ, Doppler boosting)
  - Light curve asymmetry owing to helical structure
    - Highly concentrated emission reduces this effect
  - Radio is absorbed at small scales, but could be used to trace the helical structure at larger scales

## Backup slides

### Simulations parameters

Parameter	Perucho+ 2008	Perucho+ 2010	Perucho+ 2012	Yoon+ 2016
Orbital separation	3×10 <sup>12</sup> cm	2×10 <sup>12</sup> cm	2×10 <sup>12</sup> cm	3×10 <sup>12</sup> cm
Initial jet speed	10 <sup>10</sup> cm/s	1.7×10 <sup>10</sup> cm/s	10 <sup>10</sup> cm/s	3×10 <sup>9</sup> cm/s
Wind mass loss rate	$10^{-6} \mathrm{M}_{\odot}/\mathrm{yr}$	$10^{-6} \mathrm{M}_{\odot}/\mathrm{yr}$	$10^{-6} \mathrm{M}_{\odot}/\mathrm{yr}$	$10^{-5} { m M}_{\odot}/{ m yr}$
Wind speed	2×10 <sup>8</sup> cm/s	2×10 <sup>8</sup> cm/s	2×10 <sup>8</sup> cm/s	2.5×10 <sup>8</sup> cm/s

## Clumpy wind

#### • Clumpy winds make disruption more likely to occur



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#### Perucho & Bosch-Ramon 2012

# Clumpy wind

#### Initial conditions

















### Model parameters

Parameter		Value
Stellar temperature	$T_{\star}$	$4 \times 10^4 \text{ K}$
Stellar luminosity	$L_{\star}$	$10^{39} \text{ erg s}^{-1}$
Mass-loss rate	$\dot{M}_{ m w}$	$10^{-6} \ { m M}_{\odot} \ { m yr}^{-1}$
Terminal wind speed	$v_{\infty}$	$2 \times 10^8$ cm s <sup>-1</sup>
$\beta$ -law exponent	β	0.8
Jet luminosity	$L_{i}$	$5 \times 10^{36} \text{ erg s}^{-1}$
Non-thermal energy fraction	$\eta_{ m NT}$	0.1
Acceleration efficiency	$\eta_{ m acc}$	0.1
Jet half-opening angle	$ heta_{\mathrm{j}}$	0.1 rad
Orbital separation	a	$3 \times 10^{12} \text{ cm}$
Orbital period	Т	4 days
Distance to the observer	d	3 kpc
Jet Lorentz factor	$\gamma_{i}$	1.2,3
Magnetic pressure fraction	$\eta_B^{"}$	$10^{-4}$ , $10^{-2}$ , 1
System inclination	i	$0^\circ$ , $30^\circ$ , $60^\circ$

### Energy losses



#### Particle distribution



SEDs



SEDs



i = 30°

#### Effect of the energy fraction



Molina & Bosch-Ramon 2018