

Particle acceleration in shearing flows: numerical simulation of jets and colliding winds

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Centaurus A, credit: NASA/CXC



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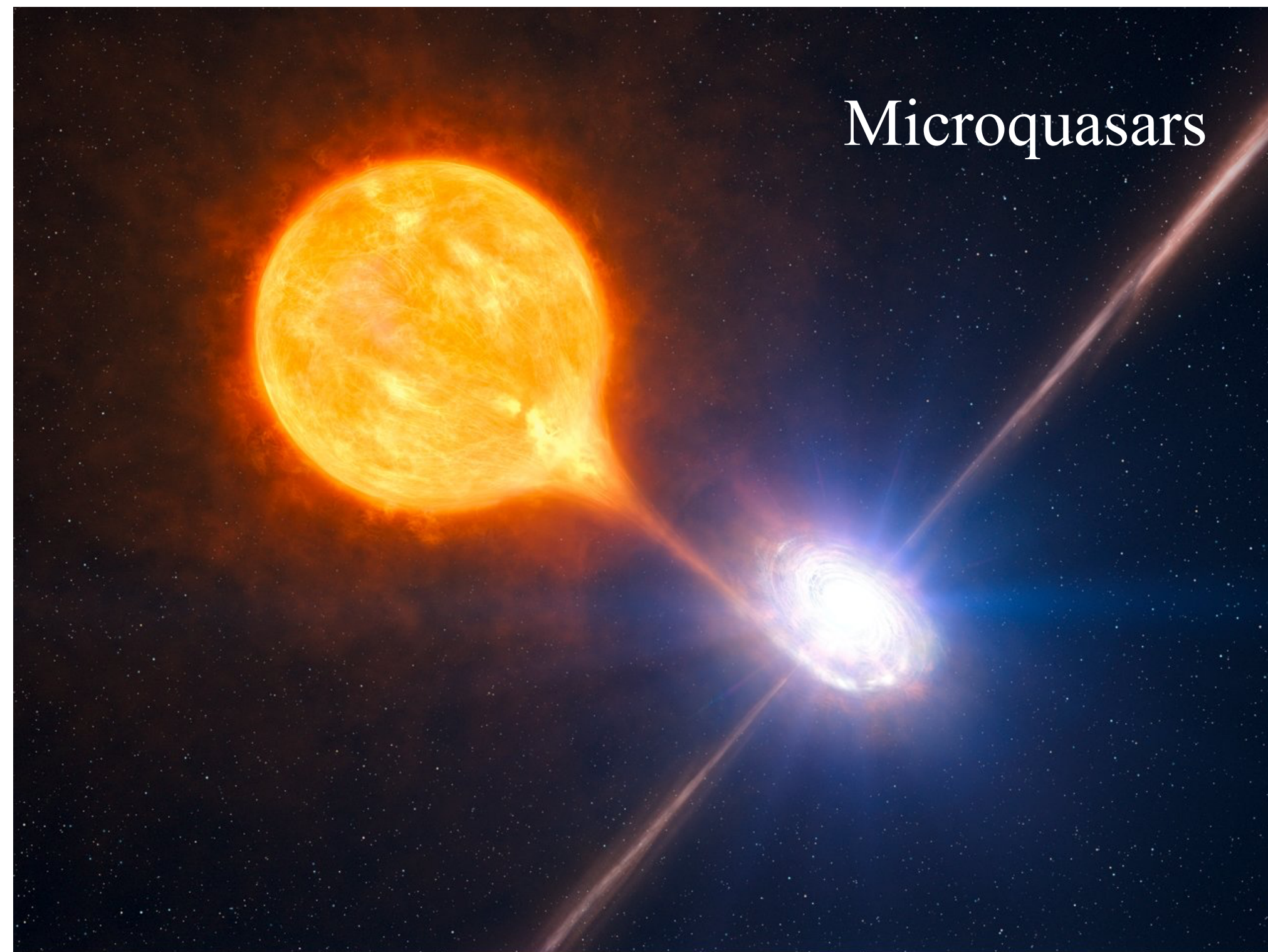
Outline

- Backgrounds and motivation
- Analytical theory of shear acceleration and its application
- RMHD and test-particle simulations of jets
- RHD simulation of colliding wind binary
- Summary

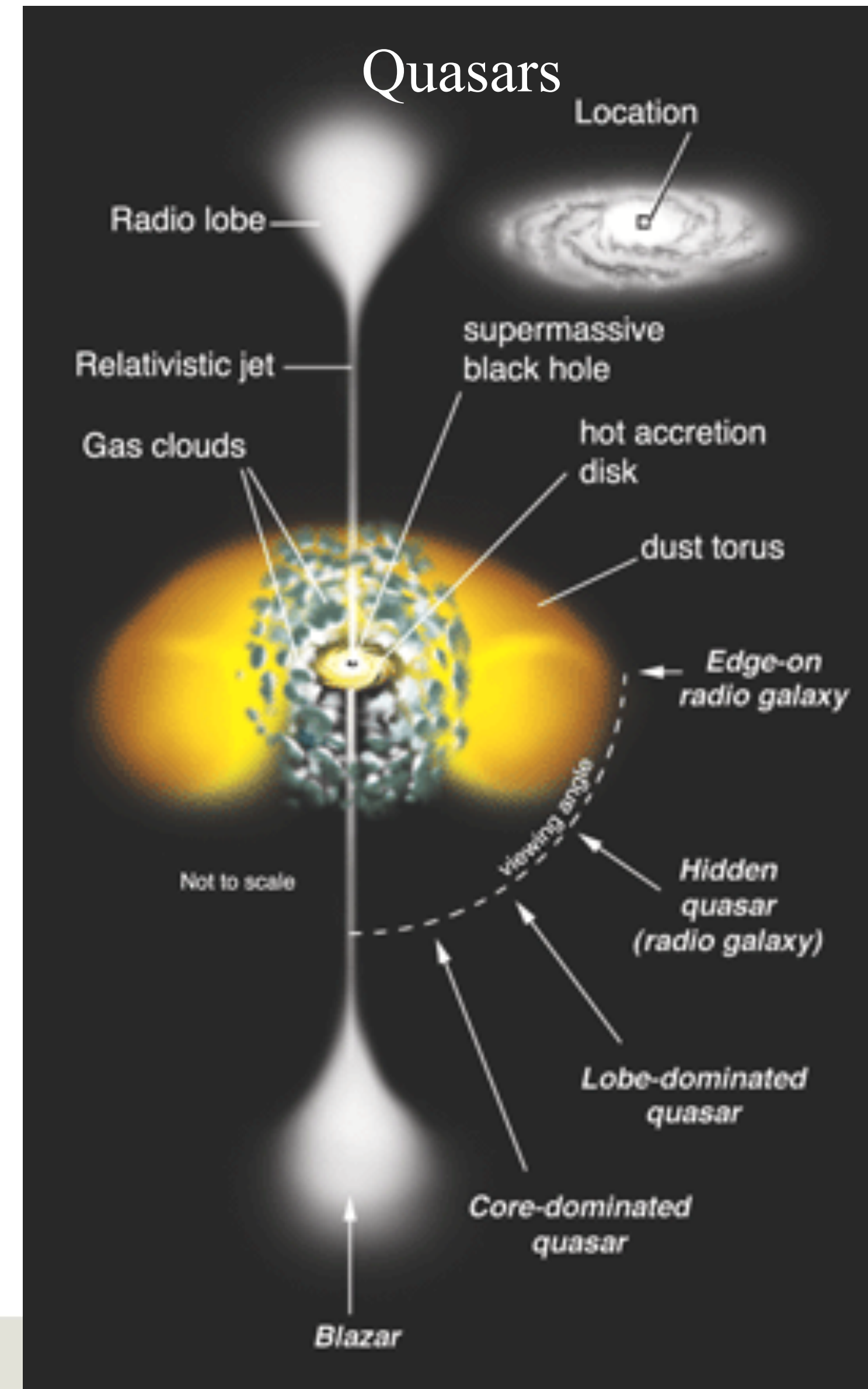


Astrophysical Jets

- Jets are ubiquitous in accreting systems



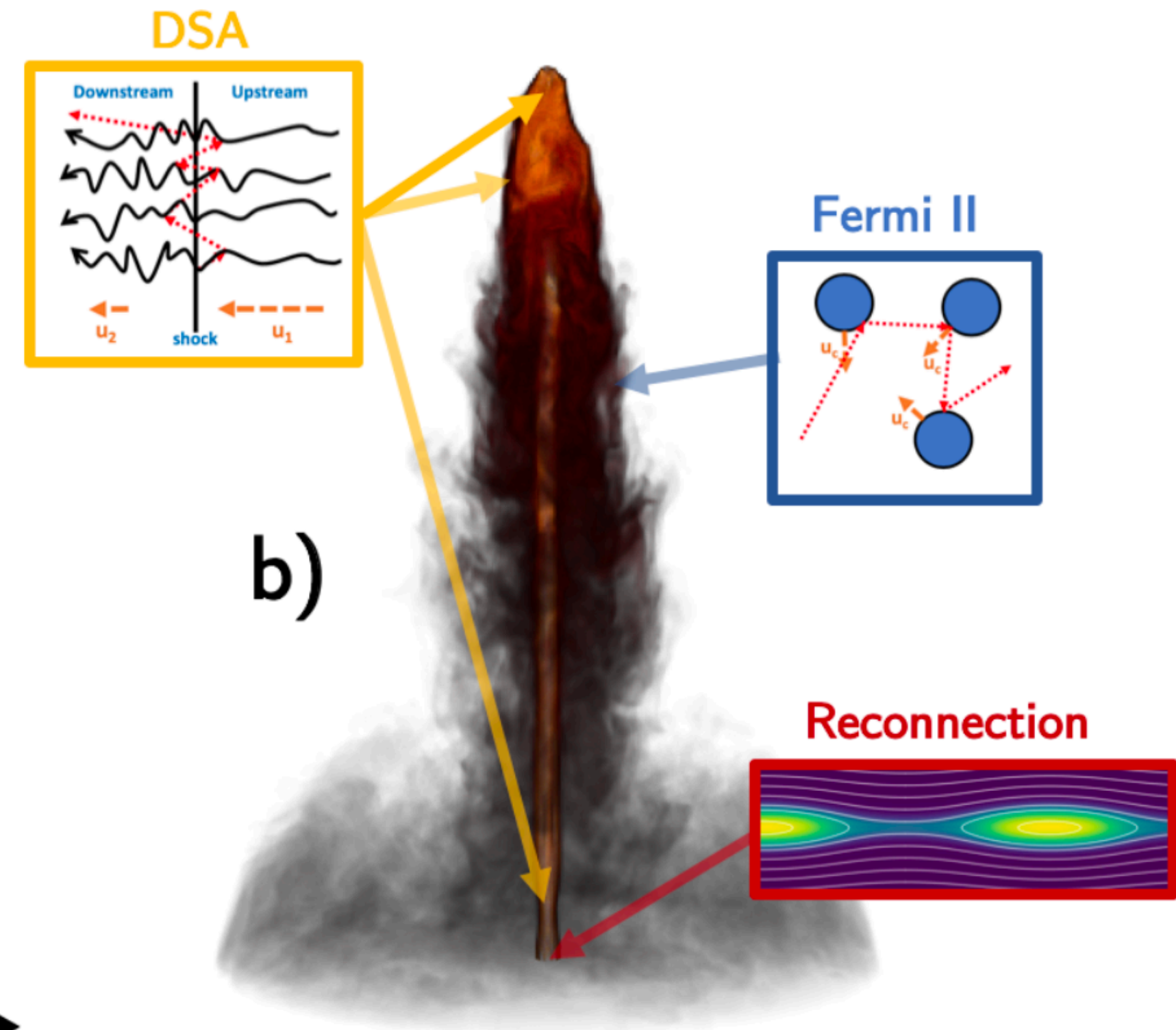
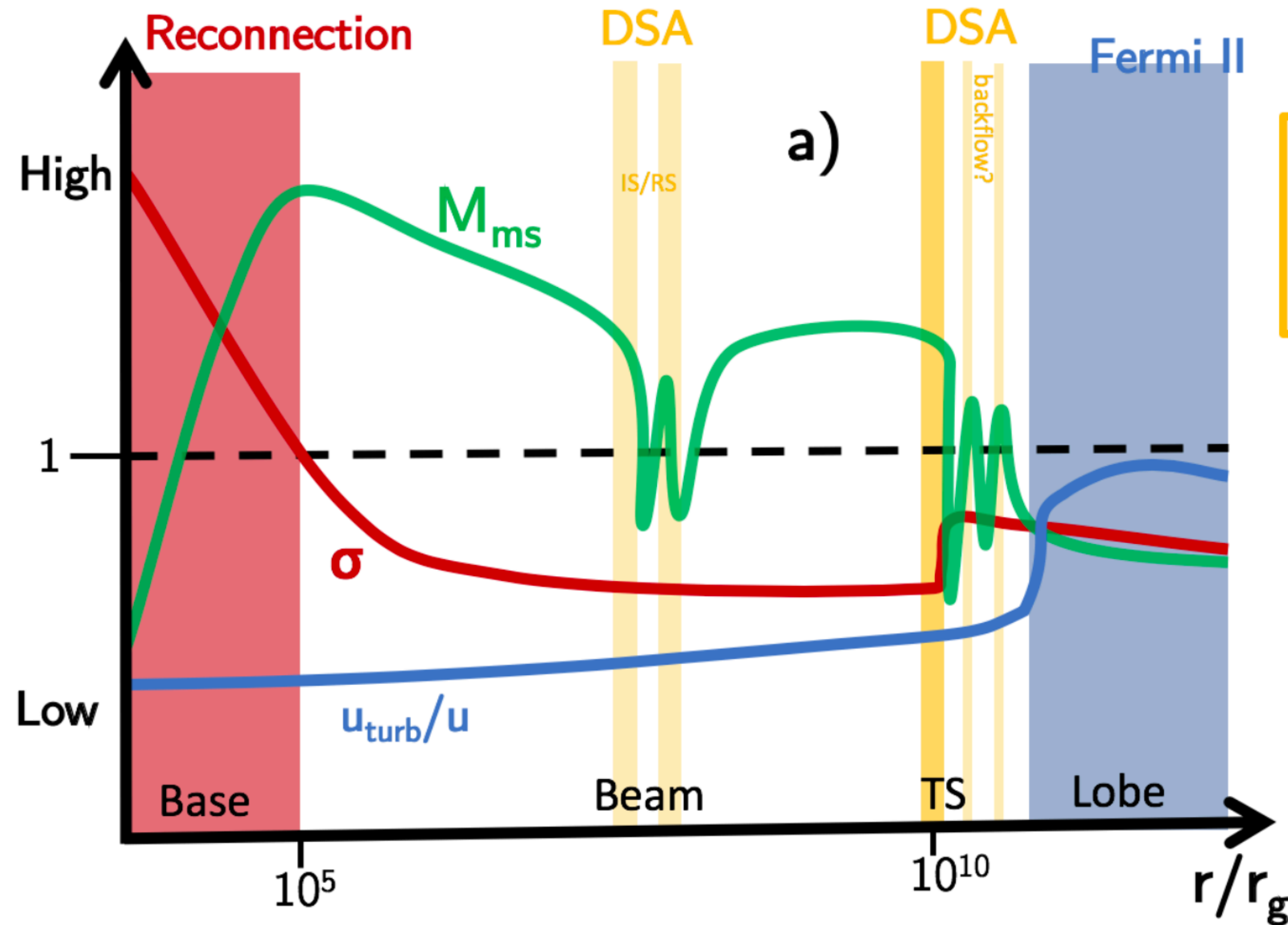
Credit: ESO/L. Calçada/M.Kornmesser



Credit: Science/
M.Kornmesser



Particle acceleration in jets

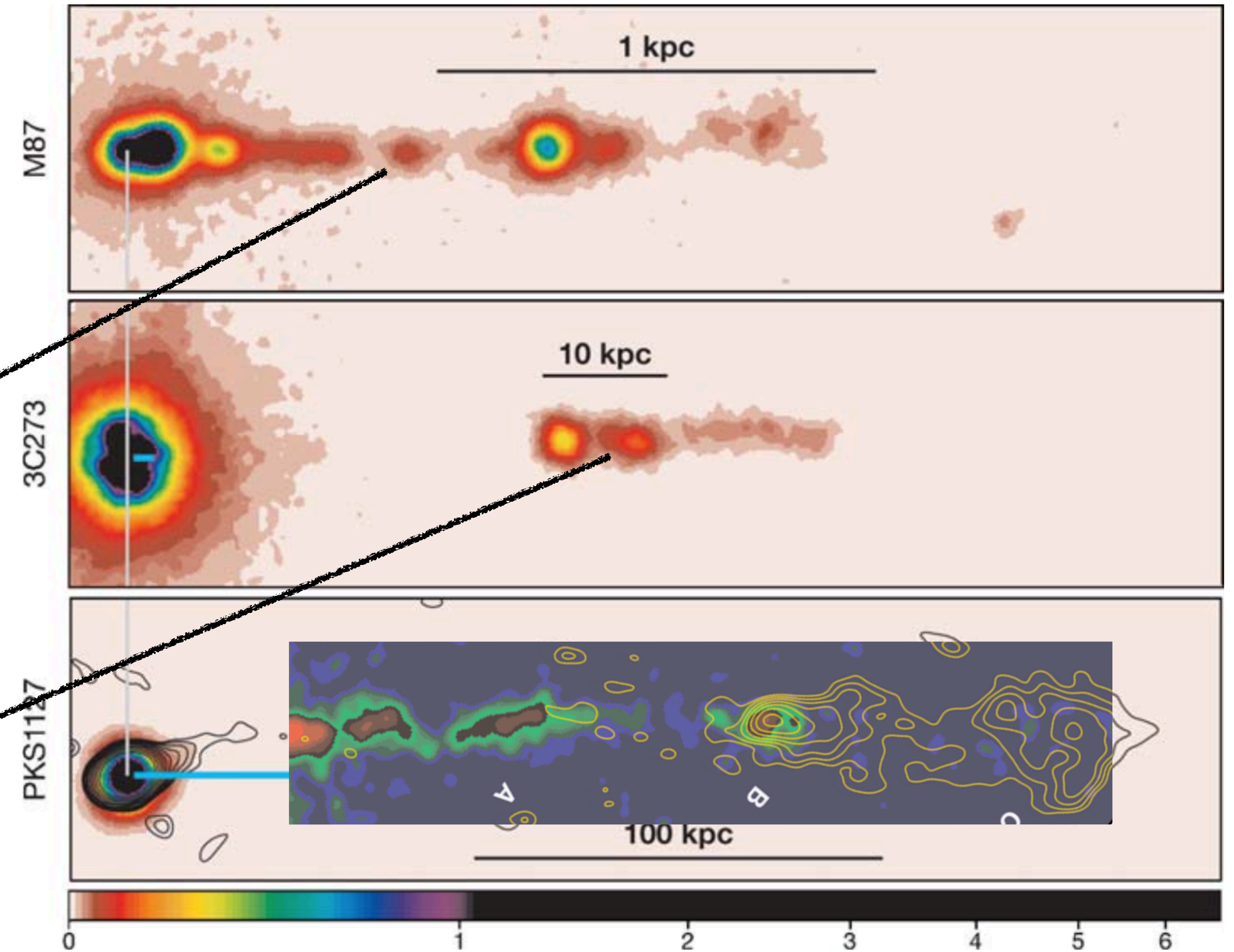
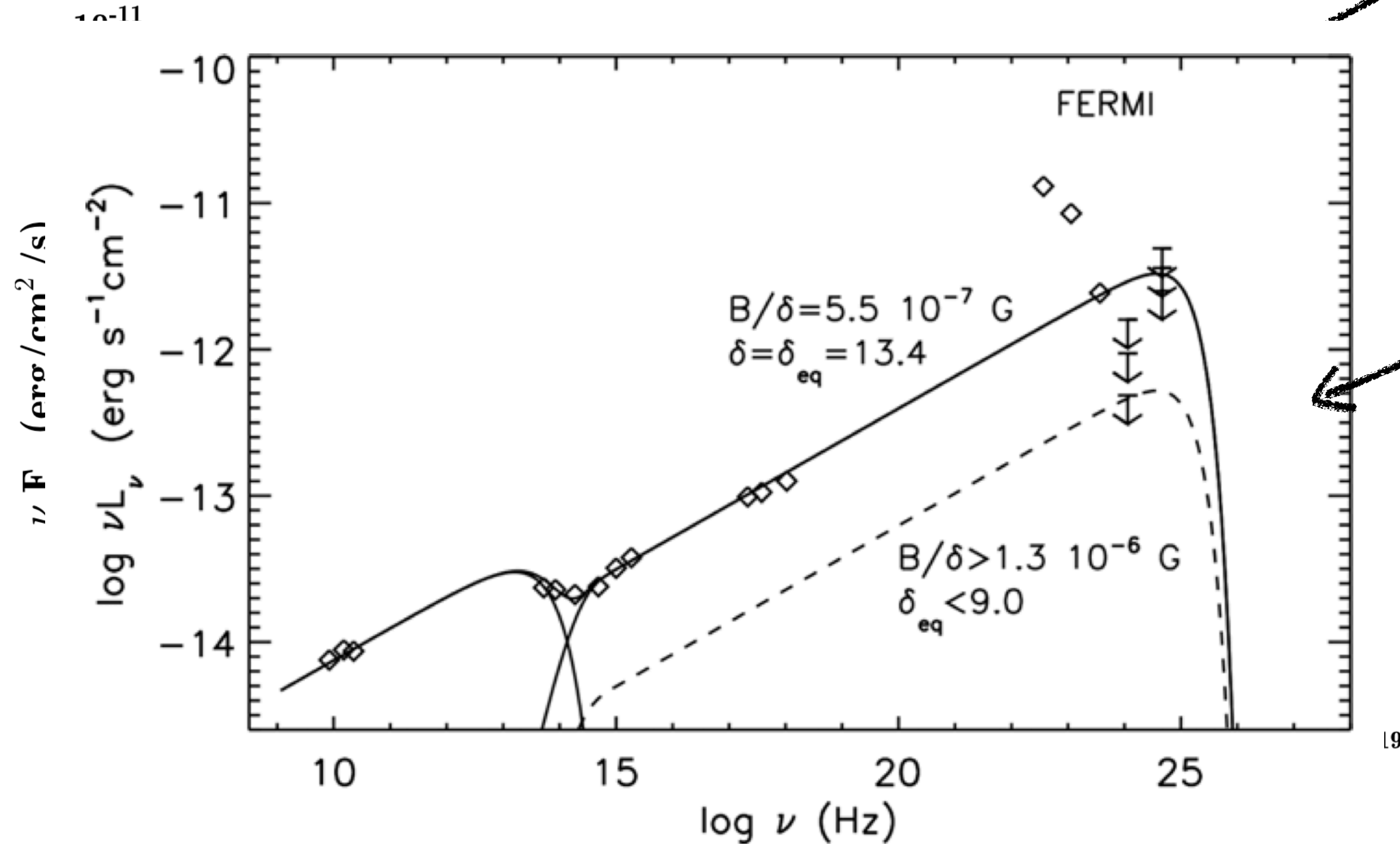


Matthews+ 2020, [arXiv:2003.06587](https://arxiv.org/abs/2003.06587)

No efficient distributed acceleration mechanism along the jet

Large-scale X-ray jets

- Quasi continuous X-rays are observed from kpc - 100kpc
- More than 100 sources @ <https://www.harvard.edu/XJET/#morph>
- Synchrotron origin is favored



Harris & Krawczynski, 2006, ARA&A, [arXiv:astro-ph/0607228](https://arxiv.org/abs/astro-ph/0607228)

Distributed acceleration required

- Synchrotron origin of X-rays requires sub-PeV electrons (PeV = 10^{15} eV):
 $E_{\text{syn}} = 2(E_e/0.1\text{PeV})^2(B/10\mu\text{G}) \text{ keV}$
- Cooling time of sub-PeV electrons:
 $\tau_c = 1.2 \times 10^3(B/10\mu\text{G})^{-2}(E_e/0.1\text{PeV})^{-1} \text{ yrs} \rightarrow \text{maximum } c\tau_c = 0.37 \text{ kpc}$
- For jet length $>$ kpc, particles accelerated by the jet head shock will cool down immediately after the shock passes (standing shocks may only exist in specific locations)
- *In-situ (re-)acceleration mechanisms are required along the jet*
 - Shear acceleration

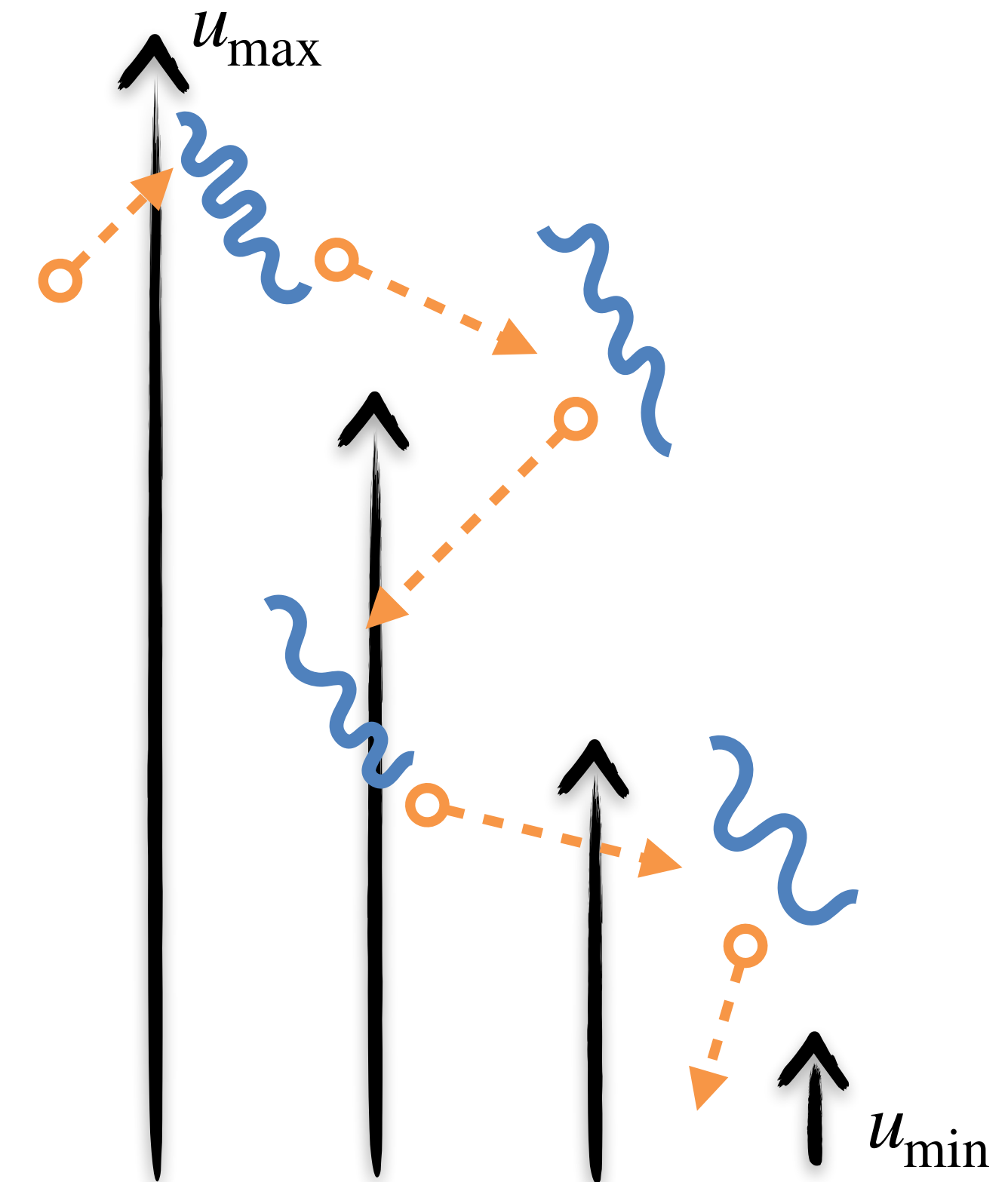
Shear acceleration

- Shear acceleration is Fermi-II type (see Rieger, 2019, [arXiv:1909.07237](https://arxiv.org/abs/1909.07237) for a review)
- Velocity profiles in jets (spine-sheath)
- Turbulences are embedded in shearing layers
- Particles scatter off turbulence (timescale: $\tau_{sc} \propto \gamma^{2-q}$, $q = 5/3$ for Kolmogorov) and sample the velocity difference

$$\frac{\langle \Delta \epsilon \rangle}{\epsilon} \propto \left(\frac{\bar{u}}{c} \right)^2 \propto \left(\frac{\partial u_z}{\partial x} \right)^2 \tau_{sc}^2$$

$$t_{\text{shear}} = \frac{\epsilon}{\Delta \epsilon} \tau_{sc} \propto \tau_{sc}^{-1} \propto \gamma^{-1/3}$$

$$\text{For comparison: } t_{\text{classical}} \propto \tau_{sc} \propto \gamma^{1/3}$$



Berezhko & Krymsky 1981; Berezhko 1982; Earl+ 1988; Webb 1989; Jokipii & Morfill 1990; Webb+ 1994; Rieger & Duffy 2004, 2006, 2016; Liu+ 2017; Webb+ 2018, 2019; Lemoine 2019; Rieger & Duffy 2019, 2021, 2022

Analytical theory: particle spectrum

- An exact solution in the steady state:

$$n(\gamma) = C_+ \gamma^{s_+} F_+(\gamma, q) + C_- \gamma^{s_-} F_-(\gamma, q)$$

$$s_{\pm} = \frac{q-1}{2} \pm \sqrt{\frac{(5-q)^2}{4} + w}$$

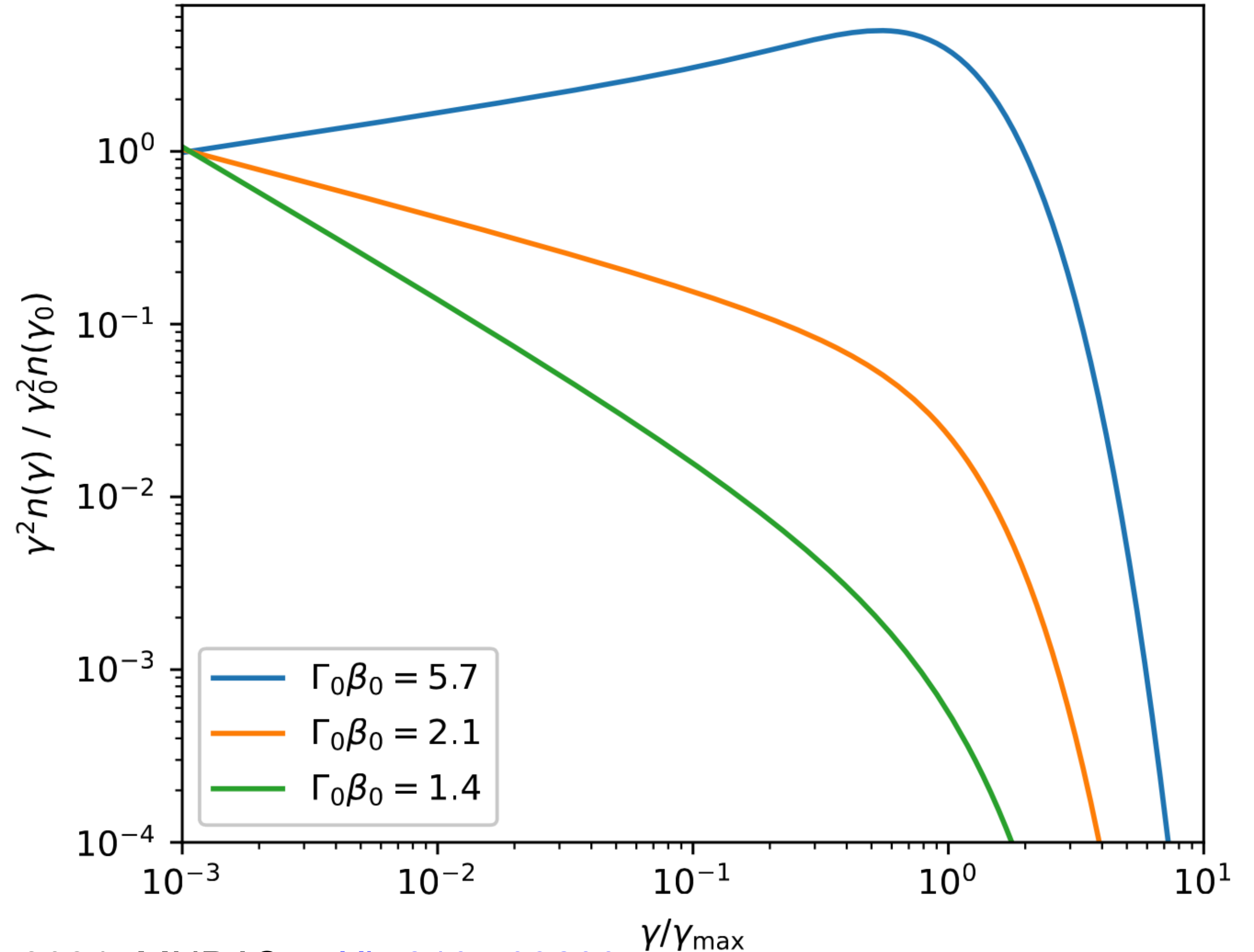
$$n \rightarrow 0 \text{ for } \gamma \rightarrow \infty$$

- Kolmogorov turbulence: $q=5/3$

- Assume a linear velocity profile

Rieger & Duffy, 2019, *ApJL*, [arXiv:1911.05348](https://arxiv.org/abs/1911.05348)

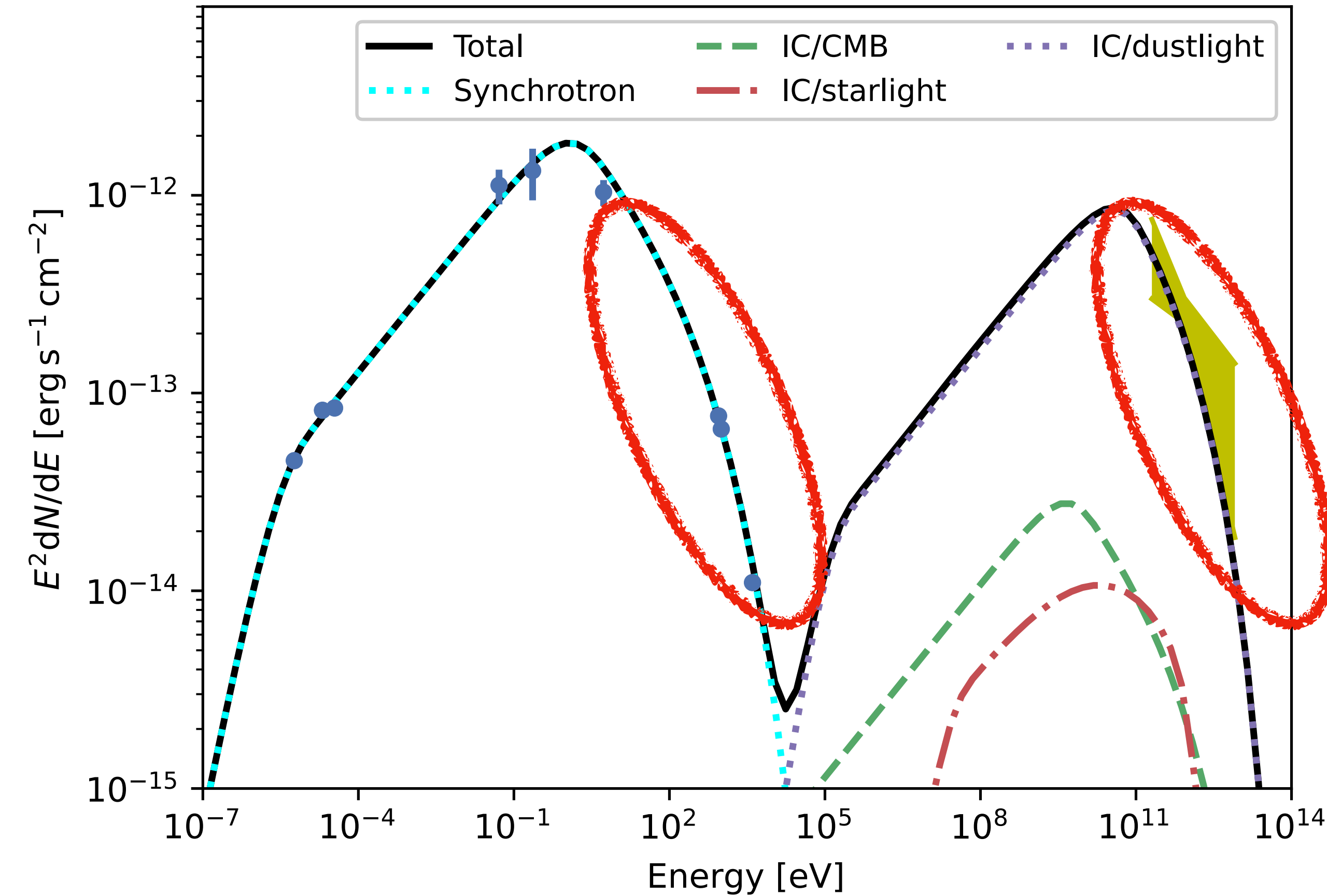
$$w = 40 \ln^{-2} \frac{(1 + \beta_0)}{(1 - \beta_0)}$$



J.S.Wang+, 2021, *MNRAS*, [arXiv:2105.08600](https://arxiv.org/abs/2105.08600)

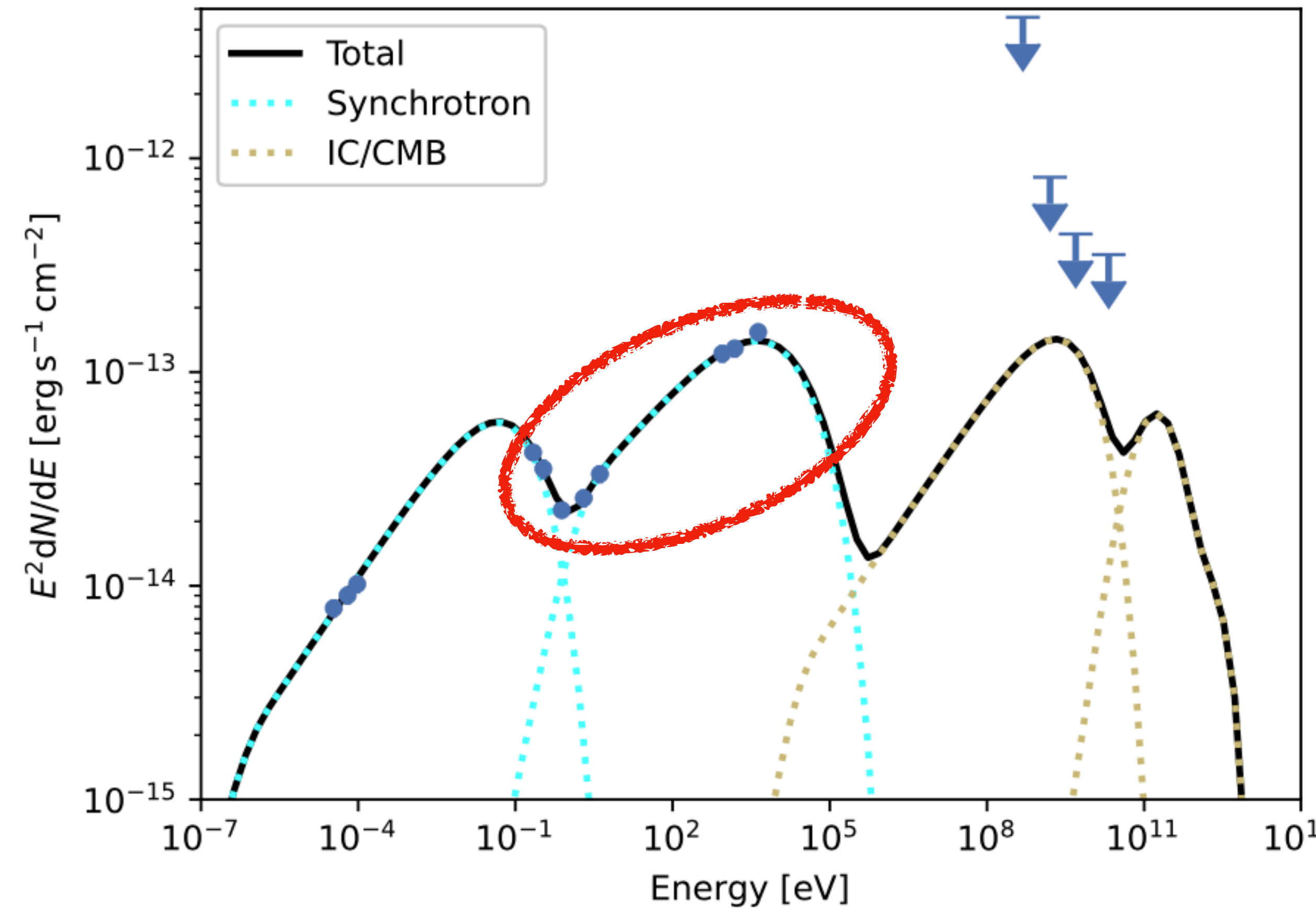
Applications

FR I: Centaurus A



IC and Syn from one population of electrons

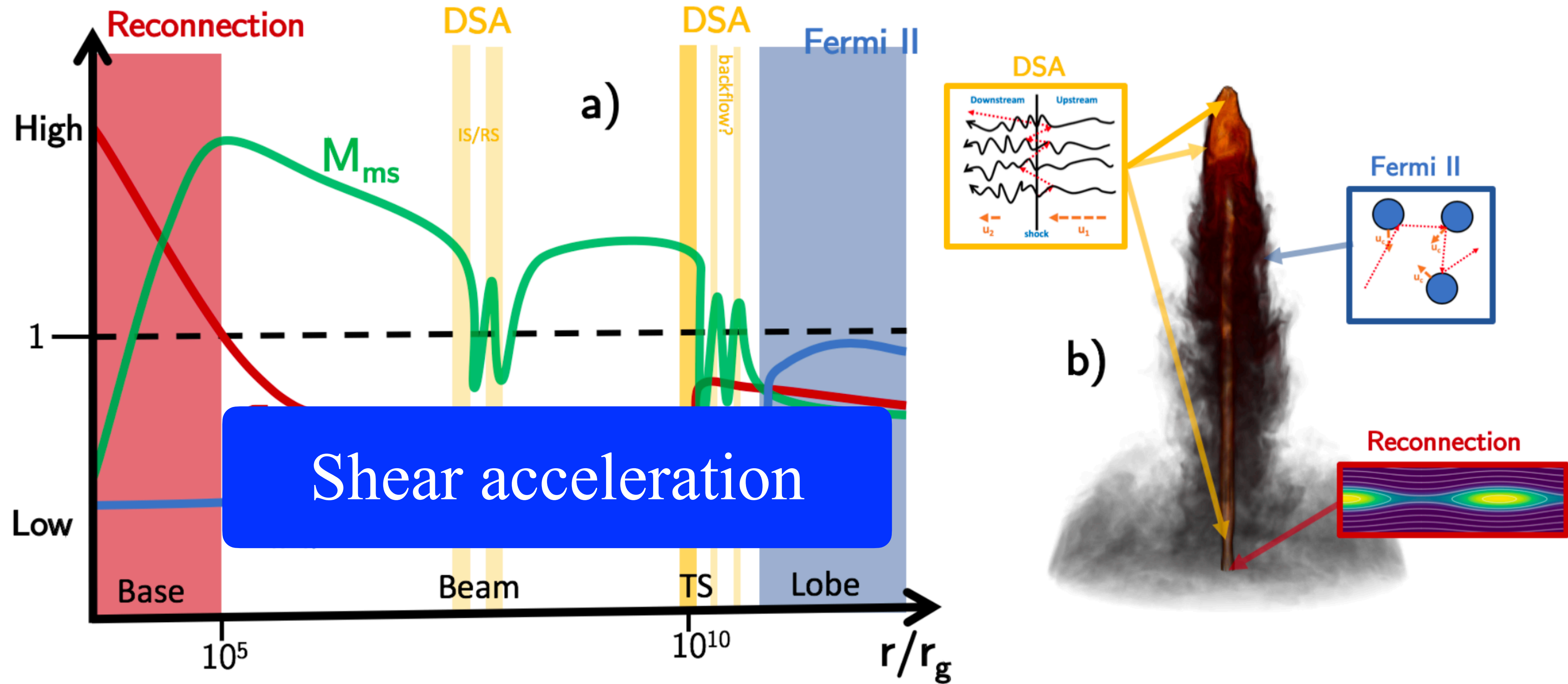
FR II: Knots A+B1 of 3C 273



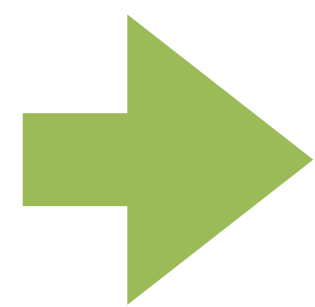
Synchrotron and IC from two populations of electrons

J.S.Wang+, 2021, *MNRAS*, [arXiv:2105.08600](https://arxiv.org/abs/2105.08600)

Particle acceleration in jets



Shear acceleration provides a good explanation to the SEDs of X-ray jets, which depends on turbulence and velocity profile



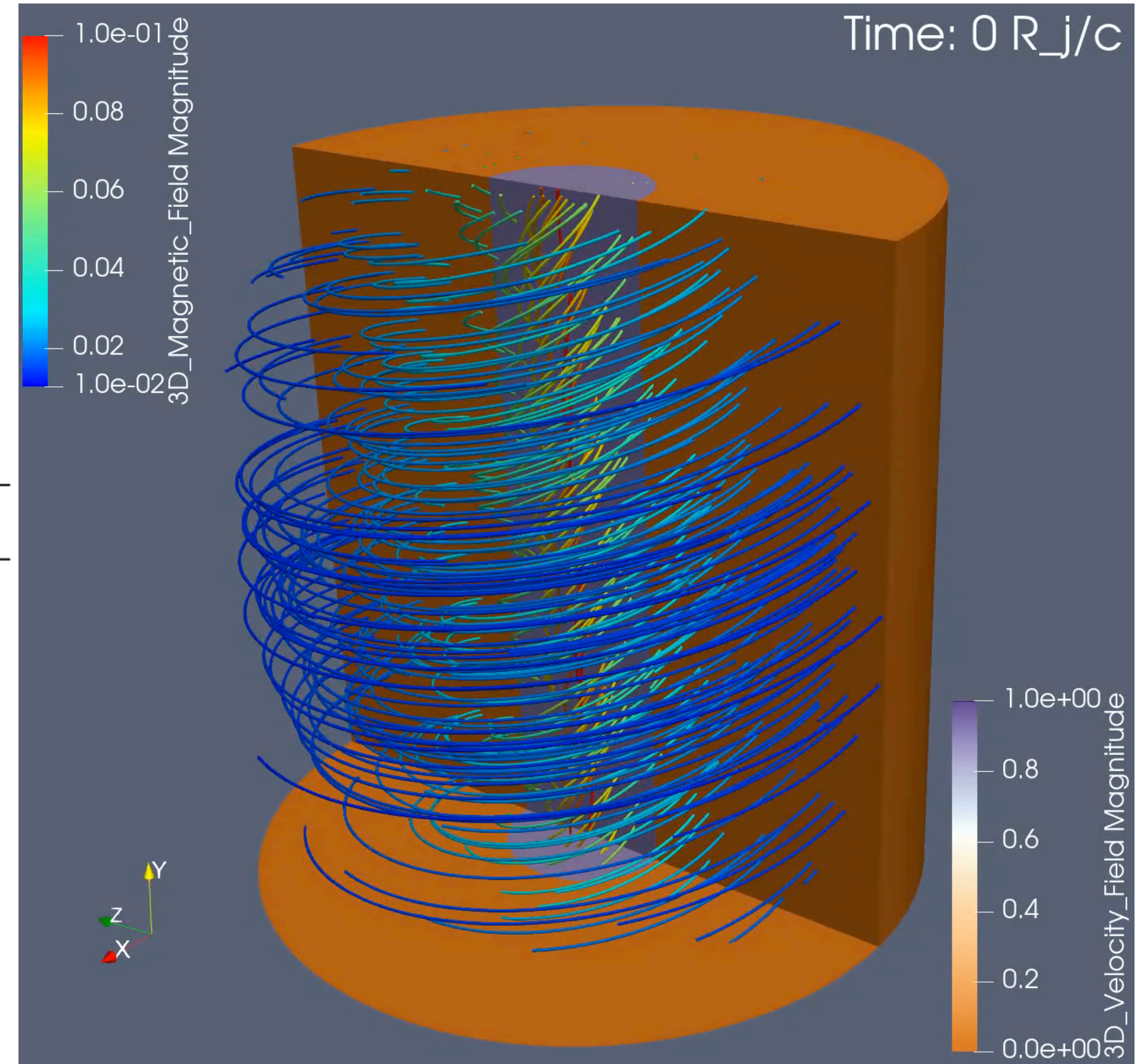
Validating shear acceleration via numeric simulations:
RMHD+Test particles

RMHD simulations with PLUTO

- Jet injected along Y axis, ambient at rest
- Periodic box along the jet axis to study the Kelvin-Helmholtz instability
- Different parameters explored
 $v \in [0.6, 0.99]$ & $\sigma \in [0.002, 0.2]$

$$\sigma_{y,\phi} = \langle B_{y,\phi}^2 \rangle / 8\pi\rho_0 c^2$$

Runs*	β_0	σ_y	σ_ϕ	Box size	Grid points	Θ_0	R_0	$L_K(\text{erg s}^{-1})$
V6B-1	0.6	10^{-1}	10^{-1}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6B-1-SB	0.6	10^{-1}	10^{-1}	$4.8R_0$	300^3	0.01	0.1kpc	1.3×10^{43}
V6B-1-LR	0.6	10^{-1}	10^{-1}	$6.0R_0$	200^3	0.01	0.1kpc	1.3×10^{43}
V6B-2	0.6	10^{-2}	10^{-2}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6BA-2	0.6	0.016	0.004	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6BT-2	0.6	0.004	0.016	$6.0R_0$	375^3	0.09	0.1kpc	1.6×10^{43}
V6B-3	0.6	10^{-3}	10^{-3}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V9B-1	0.9	10^{-1}	10^{-1}	$8.0R_0$	500^3	0.09	1 kpc	6.7×10^{45}
V9B-2	0.9	10^{-2}	10^{-2}	$8.0R_0$	500^3	0.04	1 kpc	7.0×10^{45}
V9B-3	0.9	10^{-3}	10^{-3}	$8.0R_0$	500^3	0.02	1 kpc	6.7×10^{45}
V99B-2	0.99	10^{-2}	10^{-2}	$8.0R_0$	500^3	0.07	1 kpc	7.9×10^{46}

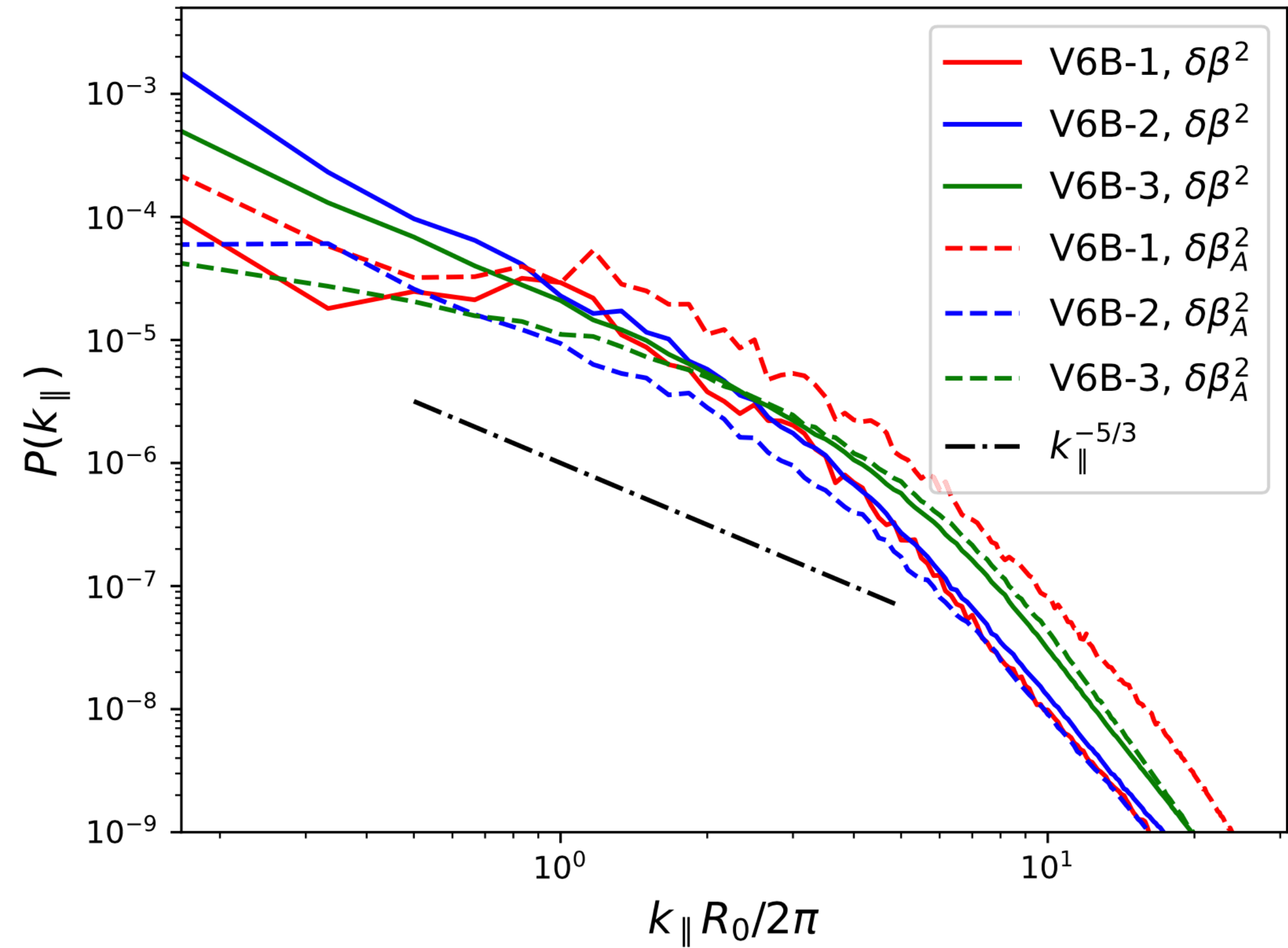
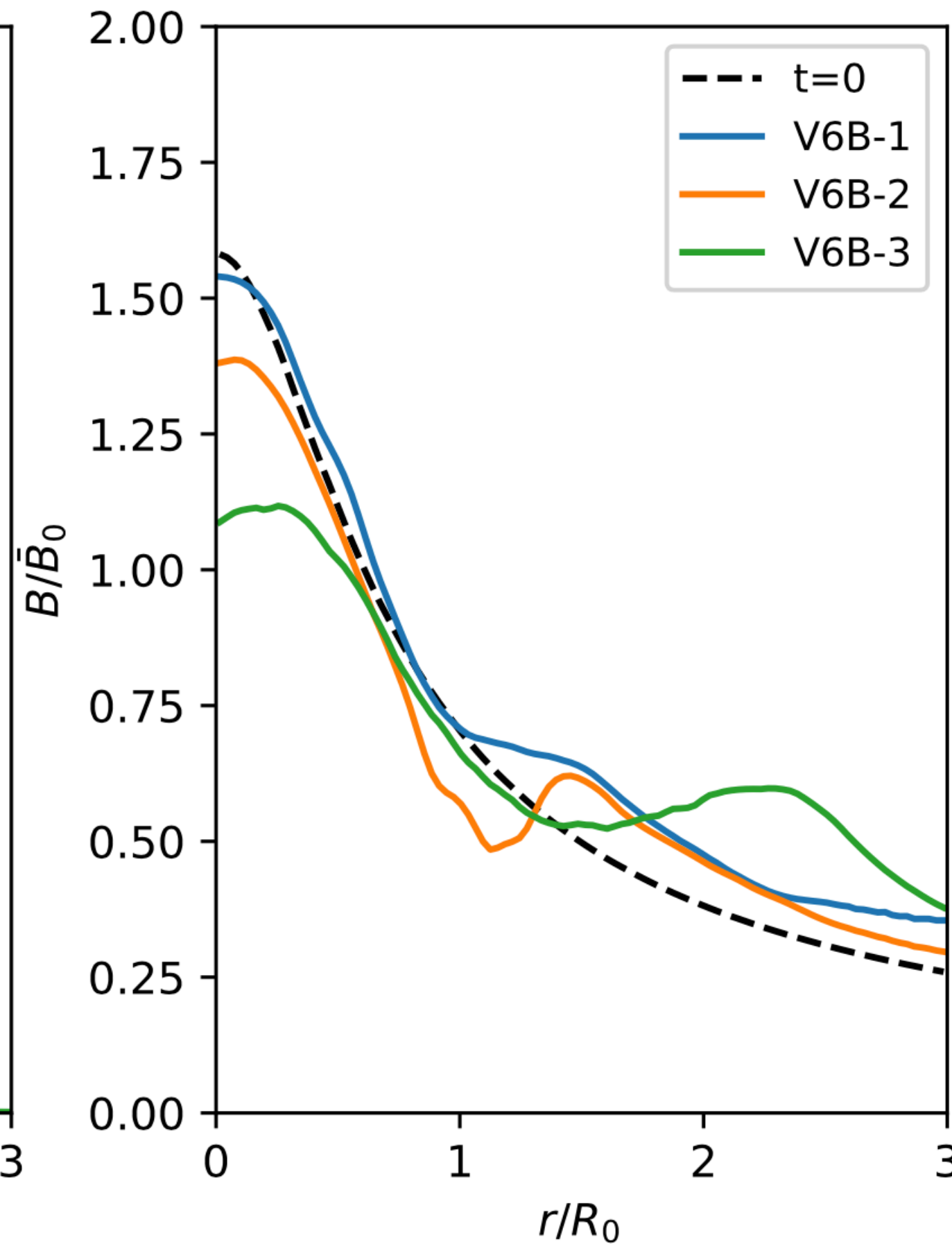
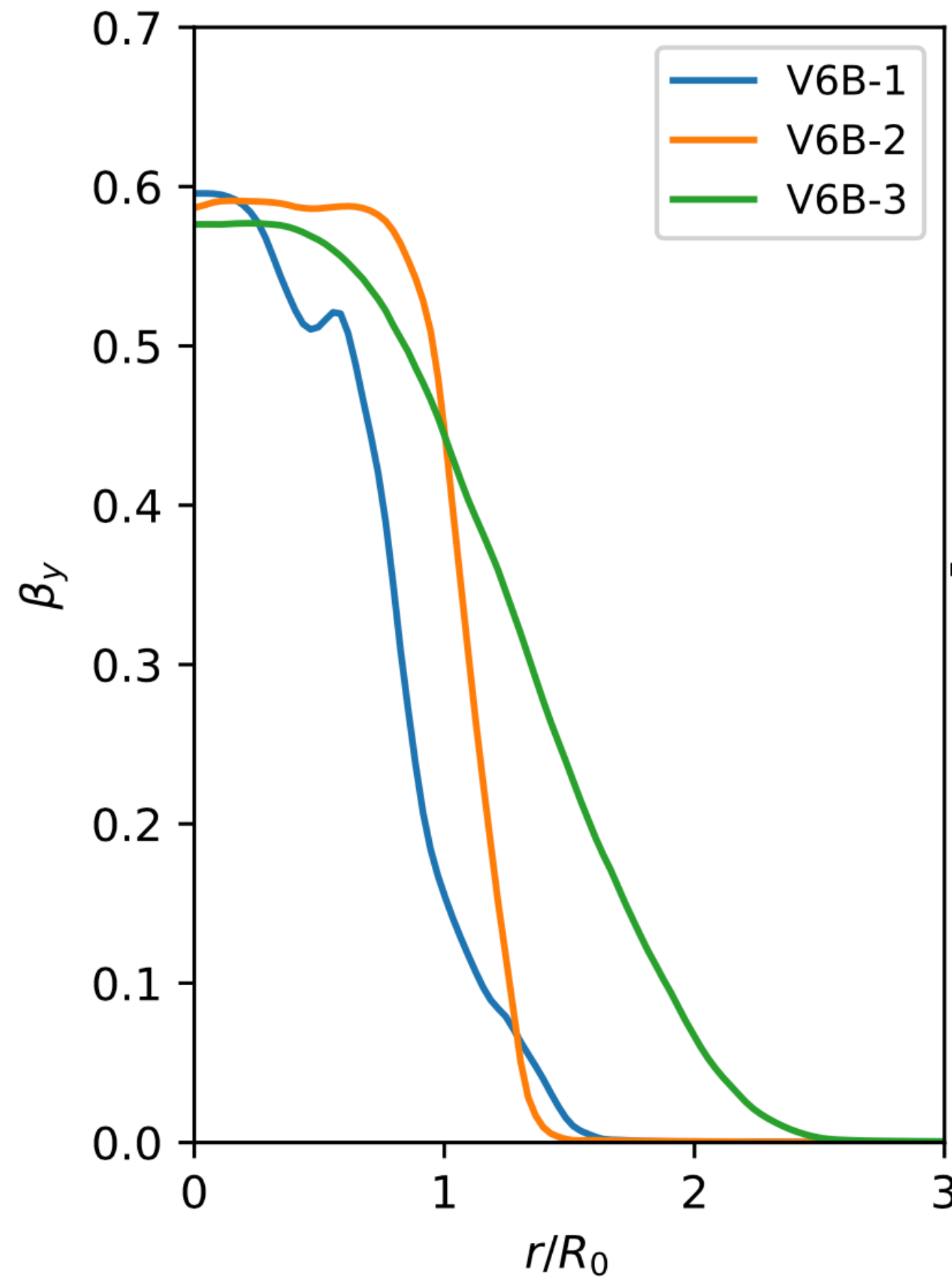


J.S.Wang+, 2023, MNRAS, [arXiv:2212.03226](https://arxiv.org/abs/2212.03226)

$v=0.6c$ cases in saturated KHI stage

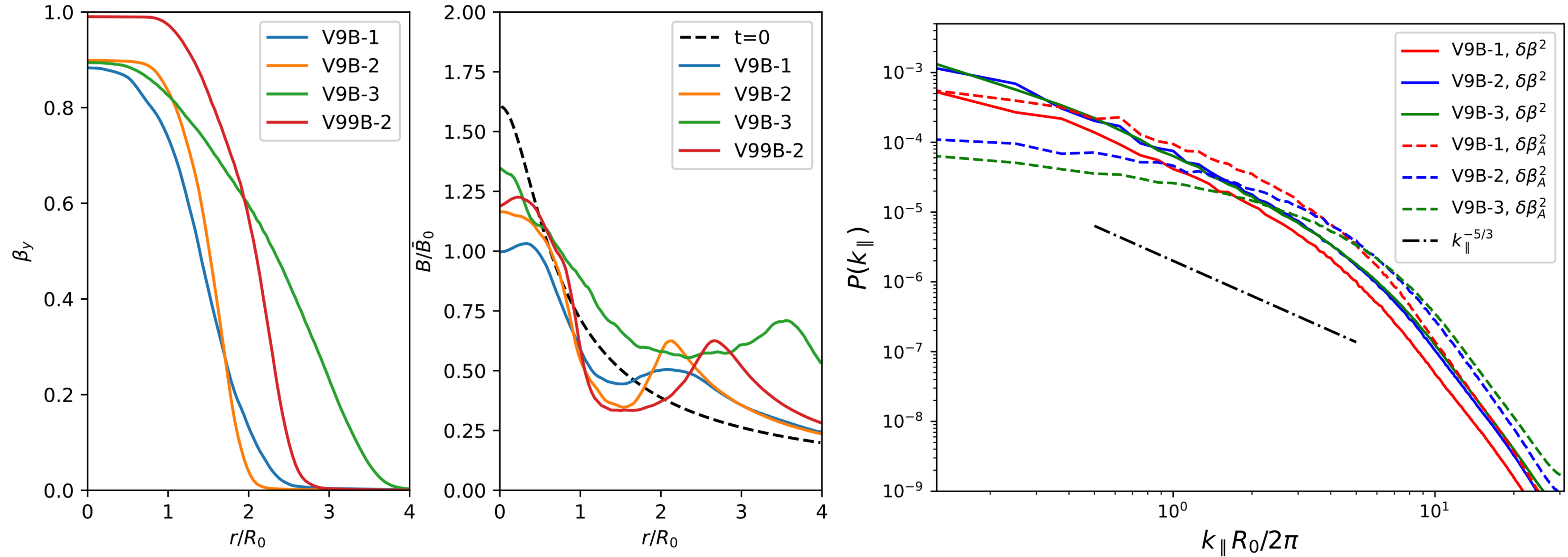
lower magnetization case: wider sheaths
& significant B field pile-up at the edge

Velocity turbulence :Kolmogorov theory



J.S.Wang+, 2023, MNRAS, [arXiv:2212.03226](https://arxiv.org/abs/2212.03226)

$v=0.9c$ cases in saturated KHI stage



J.S.Wang+, 2023, MNRAS, [arXiv:2212.03226](https://arxiv.org/abs/2212.03226)

- A sheath structure and turbulence can be self-generated via KH instability
- Higher velocities/lower magnetization lead to wider sheaths **in saturated KHI stage**
- Turbulence mostly consistent with Kolmogorov theory



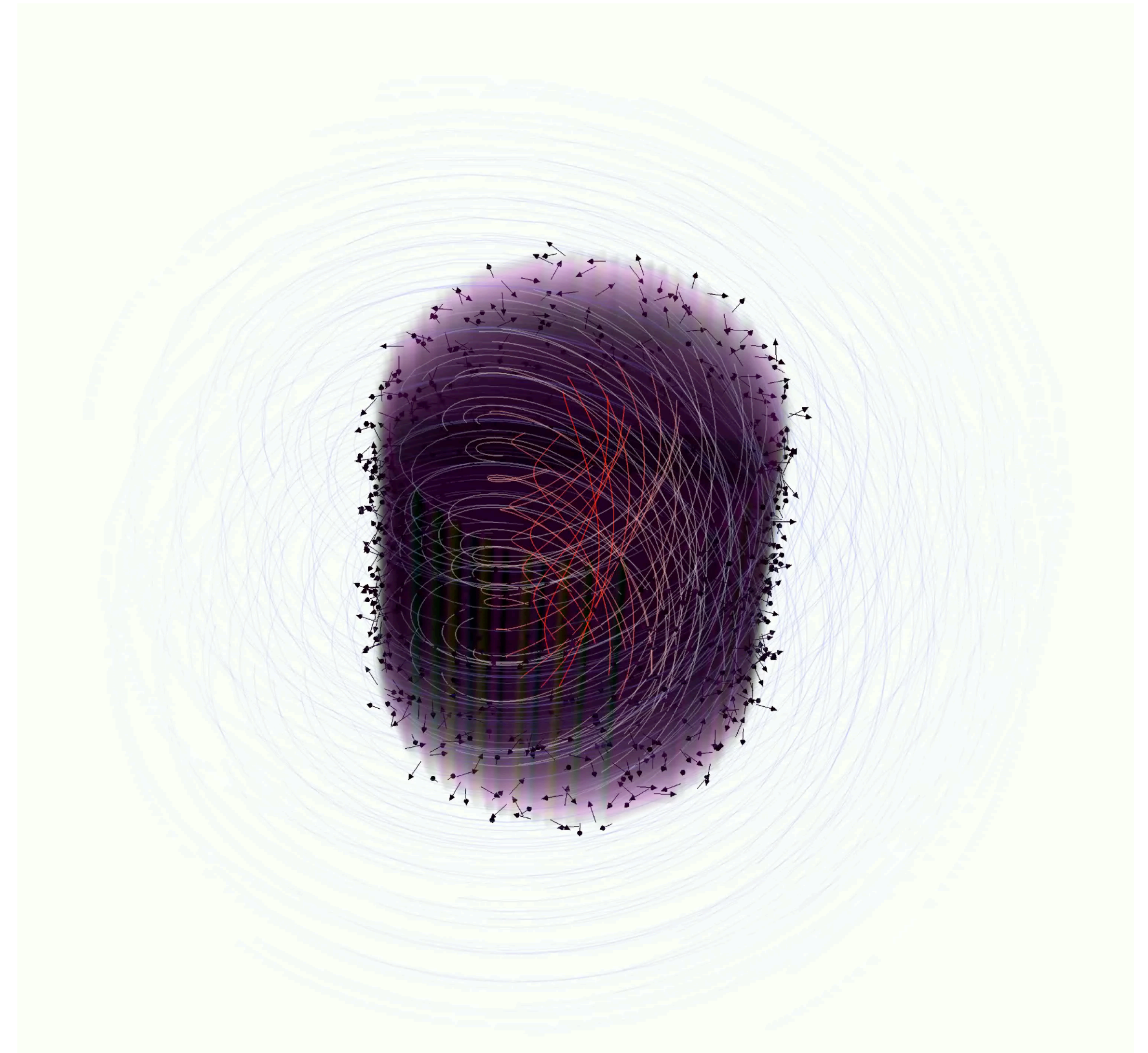
RMHD + Test-particle Simulations with PLUTO

- More self-consistent particle acceleration
- Higher-resolution runs
- Inject protons with Larmor radii at a few grid scales to avoid sub-grid physics

$$\frac{d\mathbf{x}_p}{dt} = \mathbf{v}_p$$

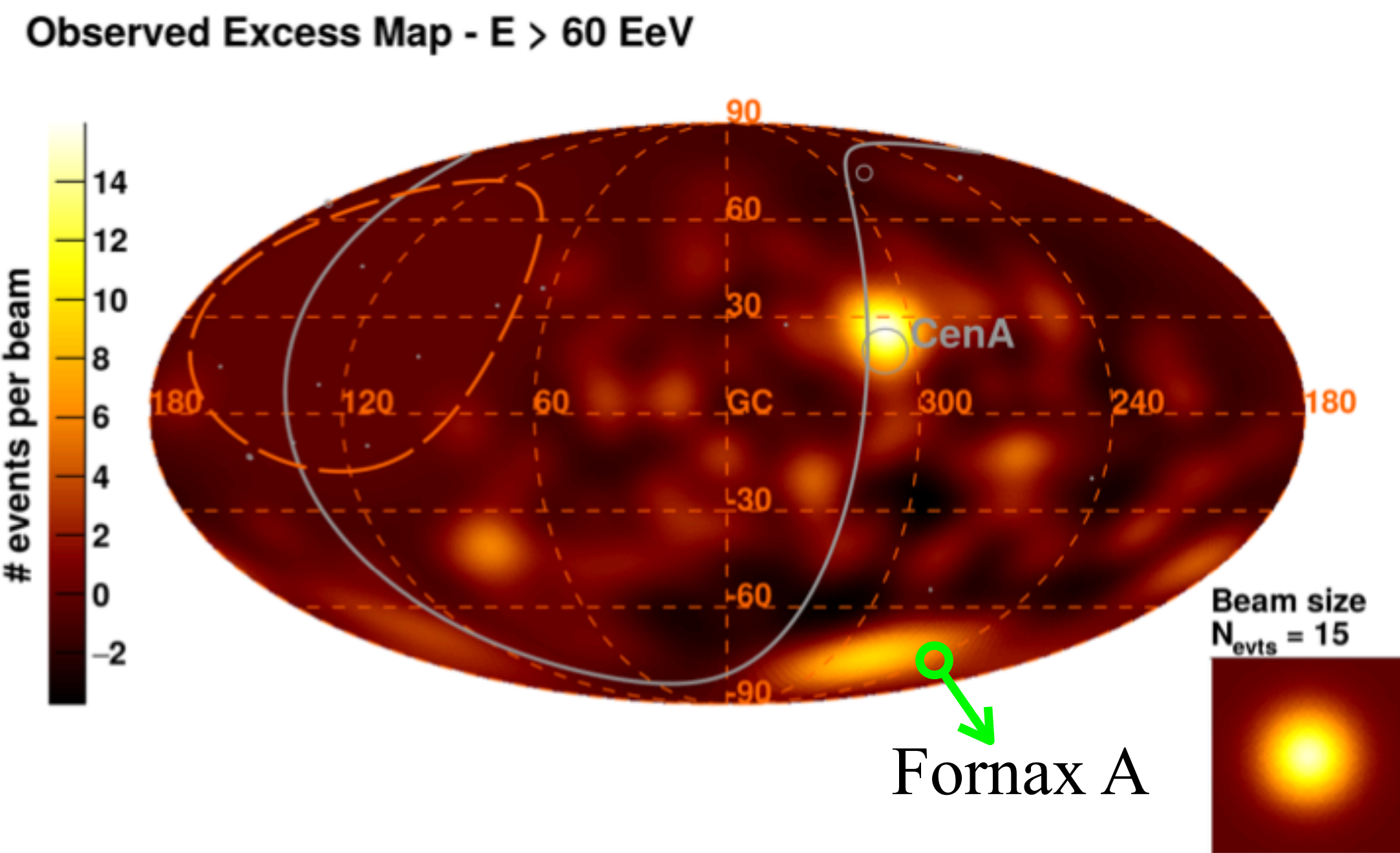
$$\frac{d(\gamma\mathbf{v})_p}{dt} = \alpha_p (c\mathbf{E} + \mathbf{v}_p \times \mathbf{B})$$

- To study the capability to accelerate UHECRs via shear acceleration

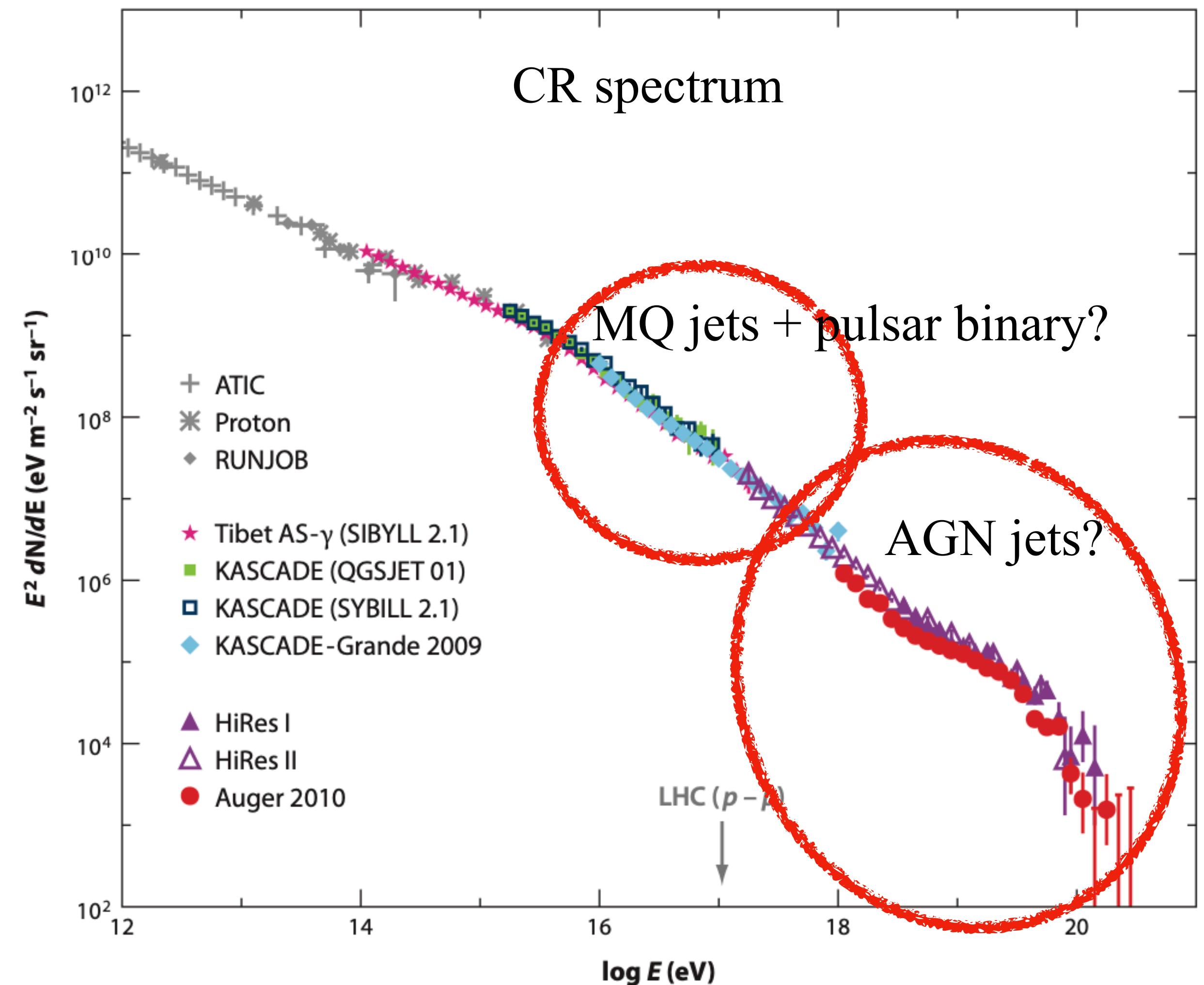


Cosmic Rays

Radio galaxies (Cen A) as UHECR sources



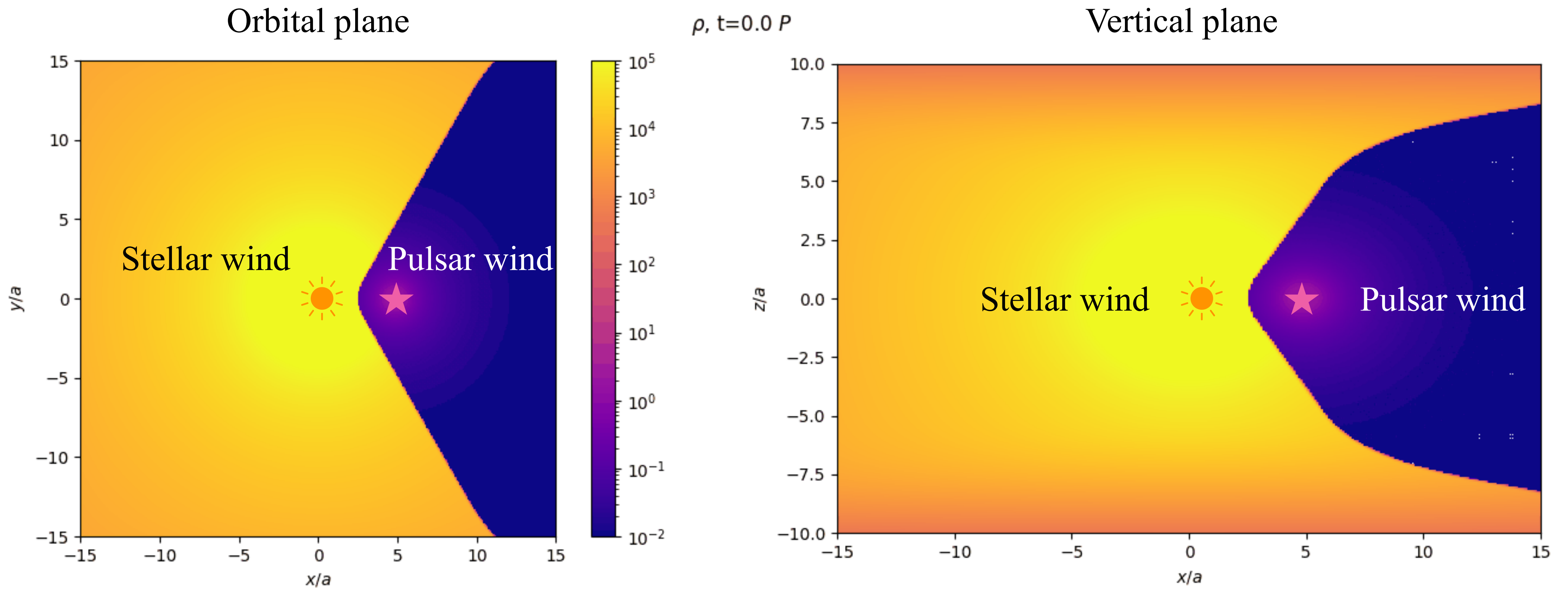
Pierre Auger Collaboration, 2018, *ApJ*, [arXiv:1801.06160](https://arxiv.org/abs/1801.06160)



Particle acceleration in pulsar binary?



Colliding wind binary: RHD simulation



Summary

- Shear acceleration is unavoidable in jets:
 - Self-generation of spine-sheath structure and turbulence via KH instability
 - Distributed particle acceleration along the jet
- Protons can achieve Hillas limit in jets via shear acceleration
 - Contribution to $>EeV$ CRs from AGN jets and >10 PeV CRs microquasar jets
- Analytical solution for shear acceleration: cut-off power-law spectra
 - Bridging theory and observations: modeling SED to get the jet parameters
- Particle acceleration in colliding pulsar wind binary to be explored (need to simplify/decouple)

